

NON-DESTRUCTIVE DETECTION FOR STRONTIUM OPTICAL LATTICE CLOCKS: TOWARDS A LATTICE CLOCK IN THE QUANTUM REGIME

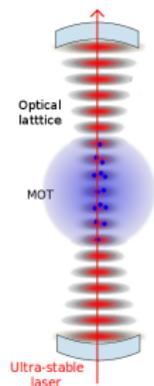
G. Vallet, S. Bilicki, E. Bookjans, R. Le Targat, and
Jérôme Lodewyck

1 OPTICAL LATTICE CLOCKS

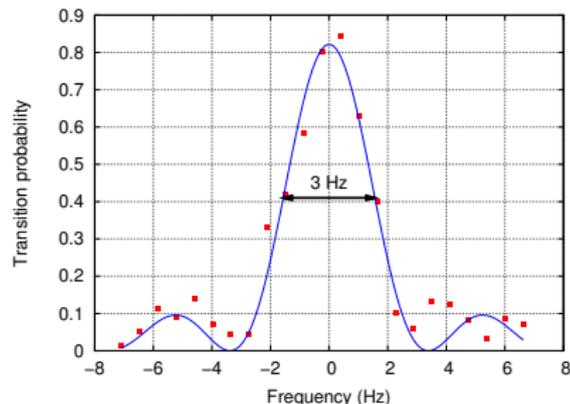
2 BEYOND THE QUANTUM PROJECTION NOISE

3 NON-DESTRUCTIVE DETECTION

OPTICAL LATTICE CLOCKS

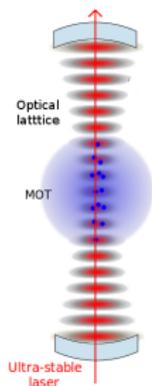


- Atoms loaded from a MOT to an **optical lattice** formed by a 1D standing wave
- Probing a narrow **optical** resonance with an ultra-stable “clock” laser
- Stabilize the clock laser on the narrow resonance



$$Q = 10^{14}$$

OPTICAL LATTICE CLOCKS



- Atoms loaded from a MOT to an **optical lattice** formed by a 1D standing wave
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COMBINE SEVERAL ADVANTAGES:

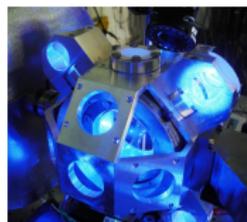
- Optical clock
- Large number of atoms (10^4)
- Lamb-Dicke regime
insensitive to motional effects
- Magic wavelength for unperturbed trapping
- Developed in many laboratories (Sr)
- \Rightarrow good candidates for a new SI second

STRONTIUM OPTICAL LATTICE CLOCKS AT SYRTE

SR1

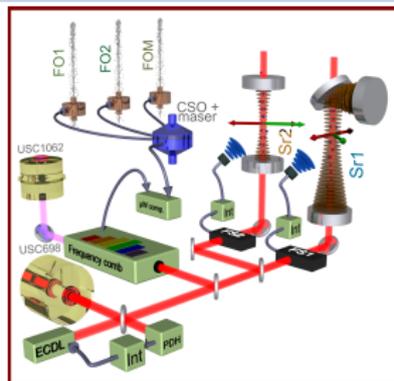


SR2

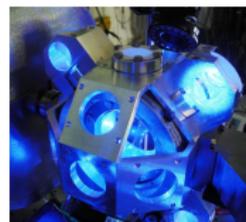


STRONTIUM OPTICAL LATTICE CLOCKS AT SYRTE

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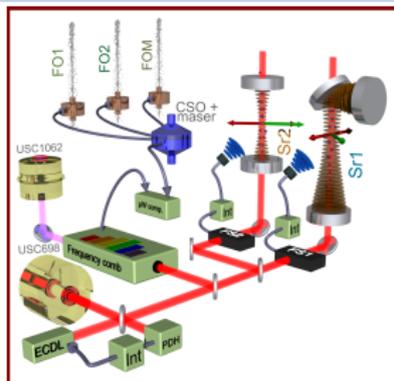


SR2

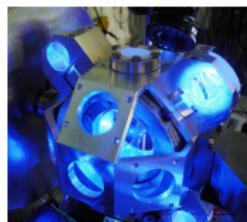


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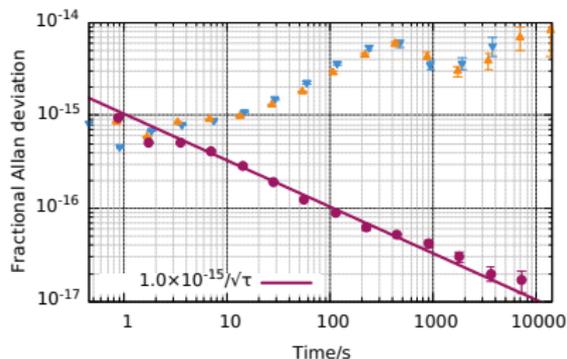
SR1



SR2



STABILITY

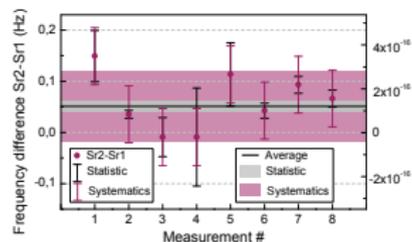


ACCURACY (in 10^{-18})

Effect	Correction	Uncertainty
Black-body radiation shift	5208	20
Quadratic Zeeman shift	1317	12
Lattice light-shift	-30	20
Lattice spectrum	0	1
Density shift	0	8
Line Pulling	0	20
Probe light-shift	0.4	0.4
AOM phase chirp	-8	8
Servo error	0	3
Static charges	0	1.5
Black-body radiation oven	0	10
Background collisions	0	8
Total	6487.4	41

RESULTS

- First agreement between two OLCs with an uncertainty beyond the accuracy of microwave clocks



$$Sr_2 - Sr_1 = 1.1 \times 10^{-16} \pm 2 \times 10^{-17}(\text{stat}) \pm 1.6 \times 10^{-16}(\text{sys})$$

Repeated agreement:

$$Sr_2 - Sr_1 = (2.3 \pm 7.1) \times 10^{-17}$$

P. Delva *et al.*, Phys. Rev. Lett. **118**, 221102 (2017)

J. Lodewyck *et al.*, Metrologia **53**, 1123 (2016)

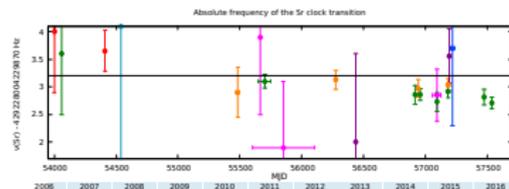
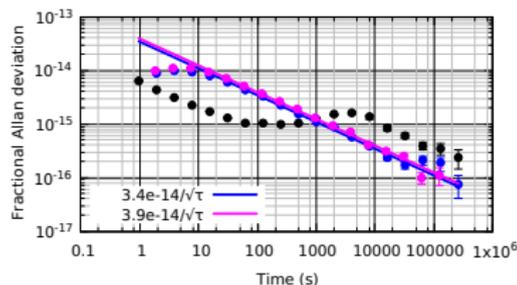
R. Tyumenev *et al.*, New Journal of Physics, **18** 113002 (2016)

C. Lisdat *et al.*, Nat. Comm. **7** 12443 (2016)

R. Le Targat *et al.* Nat. Comm. **4** 2109 (2013)

RESULTS

- First **agreement between two OLCs** with an uncertainty beyond the accuracy of microwave clocks
- Record **absolute frequency** measurement by comparing with Cs and Rb microwave fountains



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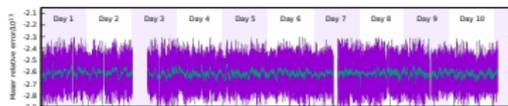
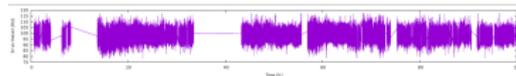
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SYRTE, PTB, JILA, Tokyo University, NICT,
NMIJ, NIM

RESULTS

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- **Continuous operation** of two Sr clocks over periods up to 3 weeks

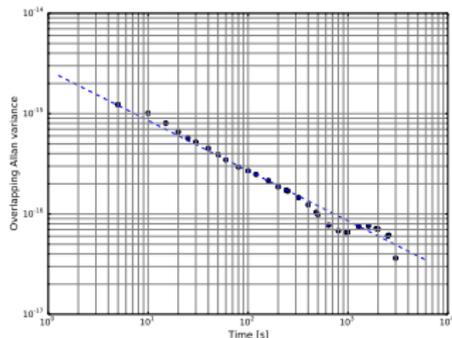


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67% to 92% uptime

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- **Optical to optical** clocks comparison with a Hg OLC



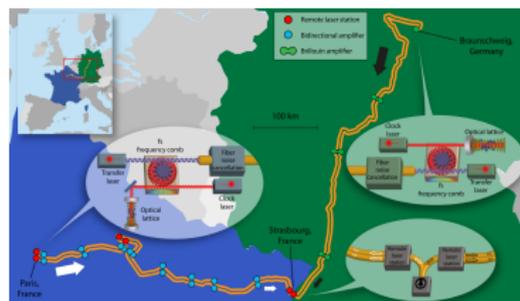
$$\text{Hg/Sr} = 2.62931420989890915 \\ \pm 5 \times 10^{-17} (\text{stat}) \pm 1.7 \times 10^{-16} (\text{sys})$$

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Best reproduced frequency ratio
(with RIKEN, Tokyo)

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- First **international comparison between optical clocks**
→ all optical comparison with phase compensated fibre links



PTB, LPL and SYRTE established a 1415 km long optical fibre link and performed in 2015 the first direct comparison of optical clocks at continental scale

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Statistical uncertainty 2×10^{-17} after $\simeq 1$ hour
150 hours of data
 $S_{rPTB}/S_{rSYRTE} - 1 = (4.7 \pm 5.0) \times 10^{-17}$

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- Bounds on **tests of Local Lorentz Invariance** with remote clock comparisons (with LPL, PTB and NPL)

correction to relativity: $|\alpha| < 10^{-8}$

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- First **international comparison between optical clocks** → all optical comparison with phase compensated fibre links
- Bounds on **tests of Local Lorentz Invariance** with remote clock comparisons (with LPL, PTB and NPL)
- First **contribution to TAI** with optical clocks

δ - duration of the TAI scale interval d .

Table 1: Estimate of d by individual PFSF measurement and corresponding uncertainties. All values are expressed in 1899-15 and are valid only for the stated period of estimation.

Standard	Period of Estimation	d	u_A	u_B	u_C	u_D	u_{TAI}	u	$u_{\text{freq Ref}(s)}$	$u_{\text{Ref}(d)}$	$u_{\text{Ref}(d)}$	Steer	Note
PTB-C11	17784 17889	-18.71	4.88	8.80	0.80	0.15	10.88	PFS/NA	7148	8.	Y	[1]	
PTB-C12	17784 17889	-0.28	3.89	11.80	0.80	0.15	12.57	PFS/NA	7148	12.	Y	[1]	
SWIT-F02	17784 17889	-1.30	0.40	0.22	0.11	0.22	0.41	PFS/NA	7381	0.24	Y	[2]	
SWIT-F08B	17784 17889	-0.91	0.20	0.27	0.11	0.22	0.49	P.F.	7328	0.30	Y	[2]	
SWIT-G02	14954 14964	0.41	0.20	0.04	0.10	0.10	0.37	0.5	[1]	0.05	N	[3]	
SWIT-G04	15179 15199	0.44	0.20	0.04	0.10	0.28	0.36	0.5	[1]	0.05	N	[3]	
SWIT-G02	17469 17479	-1.20	0.25	0.20	0.11	0.10	0.43	0.5	[1]	0.05	N	[3]	
SWIT-G04	17509 17514	-1.21	0.20	0.04	0.11	0.27	0.49	0.5	[1]	0.05	N	[3]	
SWIT-L08	15109 15104	-1.27	0.25	0.04	0.10	0.27	0.46	0.5	[1]	0.05	N	[3]	
PTB-C12F	17779 17889	-1.36	0.09	0.20	0.03	0.11	0.28	PFS/NA	7147	0.41	Y	[4]	

Notes:

[1] Continuously operating as a clock participating to TAI

[2] Report 03 MAY 2017 by IMV-SWIT

[3] Report 16 AUG 2016 by IMV-SWIT

[4] Report 02 MAR 2017 by PTB

[5] CIPM Recommendation 2 (C-2015)

[6] Update to the list of standard frequencies in Procès-Verbaux

des Séances du Comité International des Poids et Mesures, 104th meeting (2015), 2016, 41 p.

[7] Optical to microwave clock frequency ratio with a newly continuous strontium optical lattice clock, Lodewyck J.,

Blattel K., Amekong S., Hudek S.L., Shi C., Vahine G., Le Targat M., Wüchsell M., Le Coq V., Gamba J., Abgrall M.,

Rosenhahn P. and Rize S., Metrologia 53(4): 1123, 2016.

Included in Circular T 350 (Feb. 2017) as a non-steering contribution

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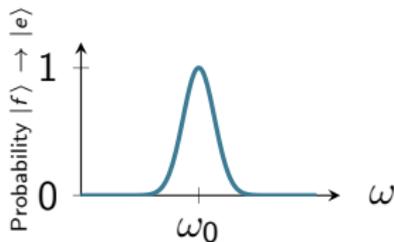
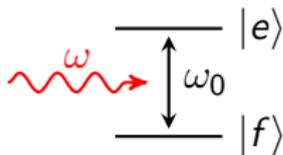
1 OPTICAL LATTICE CLOCKS

2 BEYOND THE QUANTUM PROJECTION NOISE

3 NON-DESTRUCTIVE DETECTION

QUANTUM PROJECTION NOISE: STATISTICAL MODEL

ATOMIC RESONANCE

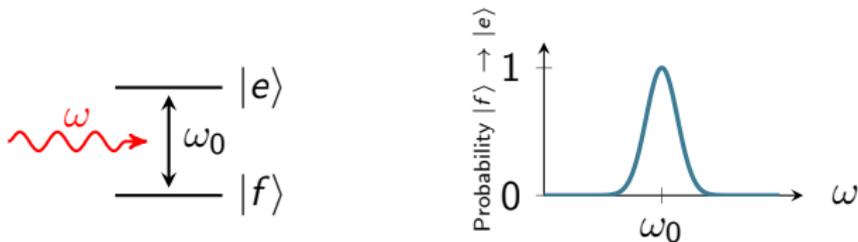


Each atom answers $|f\rangle$ or $|e\rangle$ (projection of the wave packet)

$$\Rightarrow \text{SNR} \simeq 1$$

QUANTUM PROJECTION NOISE: STATISTICAL MODEL

ATOMIC RESONANCE



Each atom answers $|f\rangle$ or $|e\rangle$ (projection of the wave packet)

$$\Rightarrow \text{SNR} \simeq 1$$

SOLUTIONS:

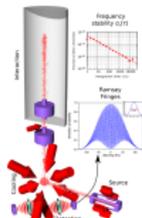
- Increase the **number of particles** $N \Rightarrow \sqrt{N}$ improvement
- Increase the **integration time** $\tau \Rightarrow \sqrt{\tau/T_c}$ improvement

Quantum projection noise limited frequency instability

$$\sigma_y(\tau) \simeq \frac{1}{\pi Q} \frac{1}{\sqrt{N}} \sqrt{\frac{T_c}{\tau}}, \quad Q = \text{quality factor}$$

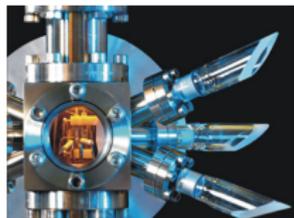
QUANTUM PROJECTION NOISE IN CLOCKS

MICROWAVE ATOMIC FOUNTAINS



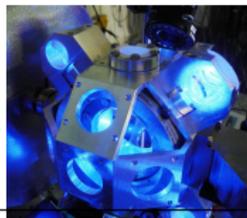
- $Q = 10^{10}$, $N = \text{a few } 10^6$
- QPN limit: $\sigma_y(\tau) = \text{a few } 10^{-14} / \sqrt{\tau}$
- Experimentally realized

OPTICAL ION CLOCKS



- $Q = \text{a few } 10^{14}$, $N = 1$
- QPN limit: $\sigma_y(\tau) = 10^{-15} / \sqrt{\tau}$
- Experimentally realized

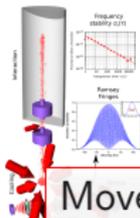
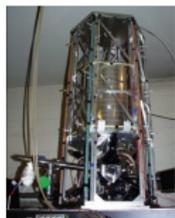
OPTICAL LATTICE CLOCKS



- $Q = \text{a few } 10^{14}$, $N = 10^4$
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- Experiments limited at $\sigma_y(\tau) = 10^{-16} / \sqrt{\tau}$
by technical noise (Dick effect)

QUANTUM PROJECTION NOISE IN CLOCKS

MICROWAVE ATOMIC FOUNTAINS

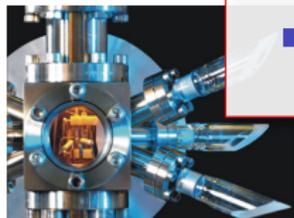


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Move to lower instabilities

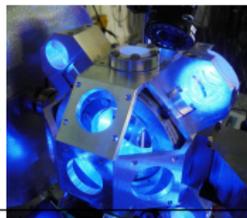
- Achieve the QPN limit in OLCs
- Develop quantum protocols to overcome the QPN

OPTICAL ION CLOCKS



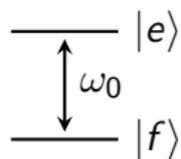
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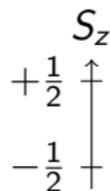


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QUANTUM PROJECTION NOISE: A SPIN MODEL



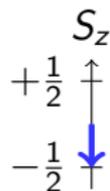
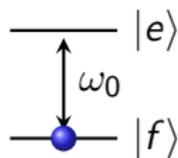
TWO LEVEL ATOM



$\frac{1}{2}$ -SPIN



QUANTUM PROJECTION NOISE: A SPIN MODEL



TWO LEVEL ATOM

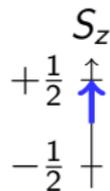
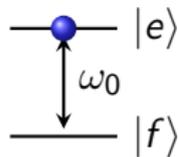
$\frac{1}{2}$ -SPIN

■ $|f\rangle$

■ $|-\rangle_z$



QUANTUM PROJECTION NOISE: A SPIN MODEL



TWO LEVEL ATOM

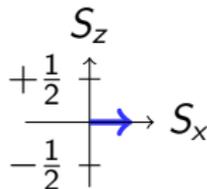
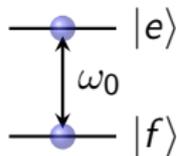
$\frac{1}{2}$ -SPIN

- $|f\rangle$
- $|e\rangle$

- $|-\rangle_z$
- $|+\rangle_z$



QUANTUM PROJECTION NOISE: A SPIN MODEL



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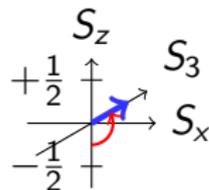
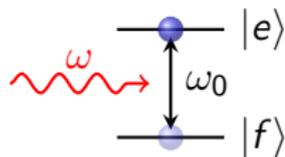
$\frac{1}{2}$ -SPIN

- $|f\rangle$
- $|e\rangle$
- $\frac{1}{\sqrt{2}}(|f\rangle + |e\rangle)$



- $|-\rangle_z$
- $|+\rangle_z$
- $\frac{1}{\sqrt{2}}(|-\rangle_z + |+\rangle_z) = |+\rangle_x$

QUANTUM PROJECTION NOISE: A SPIN MODEL



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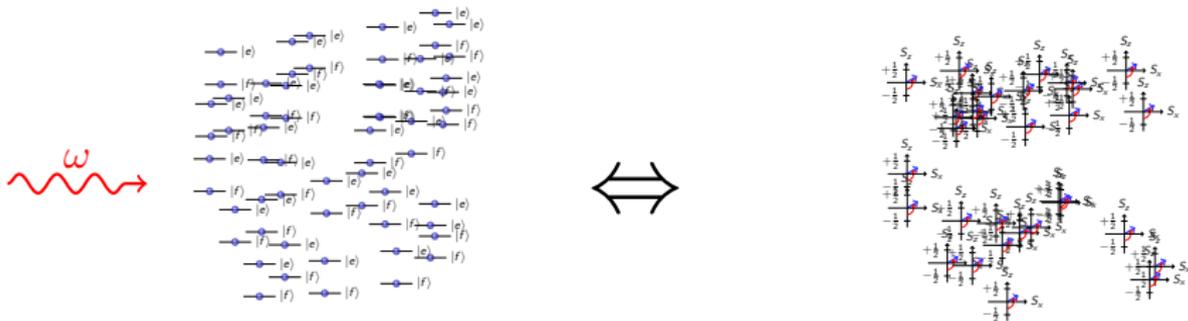
$\frac{1}{2}$ -SPIN

- $|f\rangle$
- $|e\rangle$
- $\frac{1}{\sqrt{2}}(|f\rangle + |e\rangle)$
- $\alpha|f\rangle + \beta|e\rangle, \quad |\beta|^2 = p_{|f\rangle \rightarrow |e\rangle}$



- $|-\rangle_z$
- $|+\rangle_z$
- $\frac{1}{\sqrt{2}}(|-\rangle_z + |+\rangle_z) = |+\rangle_x$
- $|+\rangle_3$

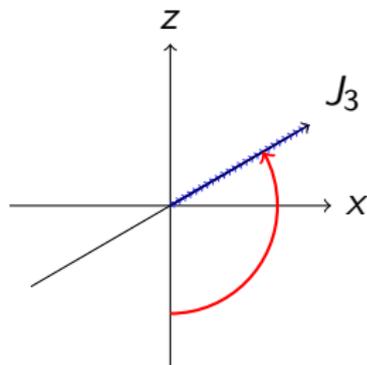
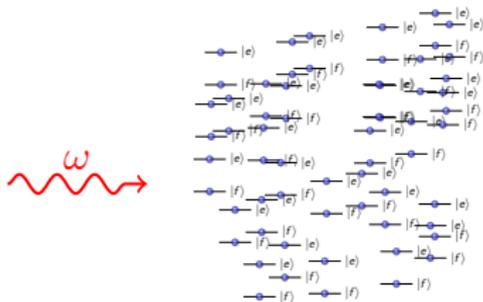
QUANTUM PROJECTION NOISE: A SPIN MODEL



N TWO LEVEL ATOMS

N $\frac{1}{2}$ -SPINS

QUANTUM PROJECTION NOISE: A SPIN MODEL

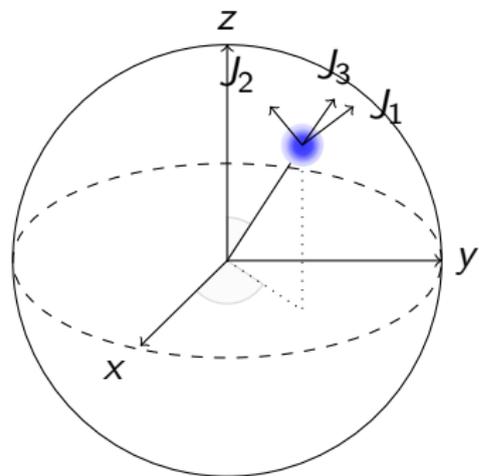


N TWO LEVEL ATOMS

N $\frac{1}{2}$ -SPINS

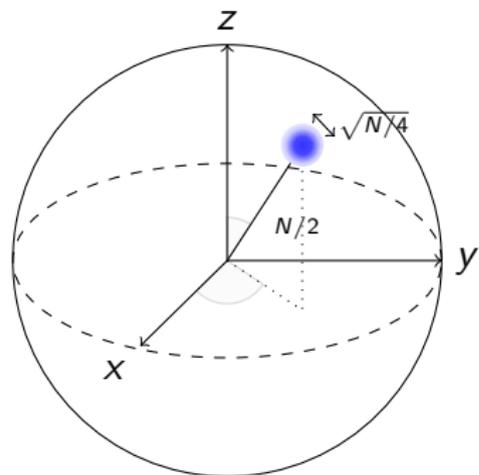
- $\vec{J} = \sum_{i=1}^N \vec{S}^{(i)}$
- $J_3 \simeq \frac{N}{2}$

SPIN SQUEEZING



- $[J_1, J_2] = iJ_3 \simeq i\frac{N}{2}$
- $\Rightarrow \Delta J_1 \Delta J_2 \geq \frac{N}{4}$
- Similar to $[X, P] = i\hbar$ systems
 - Quantum harmonic oscillator
 - Quantum optics

SPIN SQUEEZING

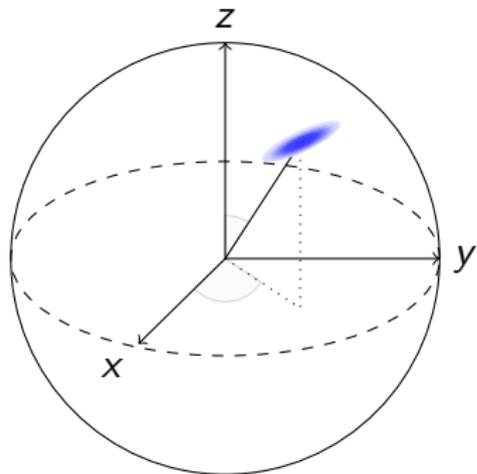


- $[J_1, J_2] = iJ_3 \simeq i\frac{N}{2}$
- $\Rightarrow \Delta J_1 \Delta J_2 \geq \frac{N}{4}$
- Similar to $[X, P] = i\hbar$ systems
 - Quantum harmonic oscillator
 - Quantum optics

QUANTUM PROJECTION NOISE

- **Symmetric** uncertainty area
- $\Delta J_1 = \Delta J_2 = \sqrt{\frac{N}{4}}$
- $\text{SNR} = \frac{N/2}{\sqrt{N/4}} = \sqrt{N}$ (QPN)

SPIN SQUEEZING



- $[J_1, J_2] = iJ_3 \simeq i\frac{N}{2}$
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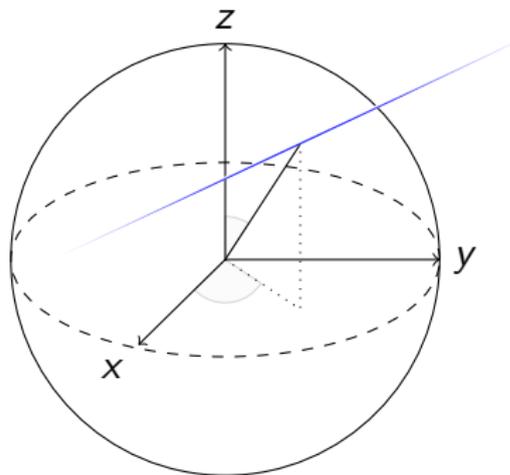
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SPIN SQUEEZING

- **Asymmetric** uncertainty area
- $\Delta J_2 < \Delta J_1$
- Limit: $\Delta J_1 \simeq N \rightarrow \Delta J_2 \simeq 1$
 $\Rightarrow \text{SNR} \simeq N$
(Heisenberg limit)

SPIN-SQUEEZING: REQUIREMENTS

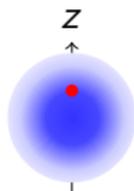
HOW TO ACHIEVE SPIN SQUEEZING

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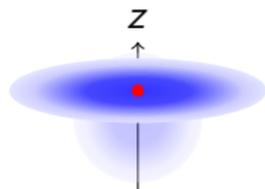


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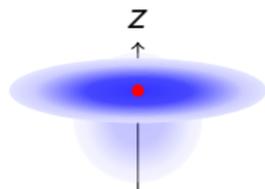


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- 2 Subsequent measurement correlated

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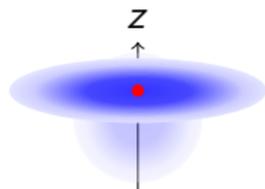


- 1 QPN limited measurement
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- 3 Design a protocol to achieve a sub-QPN resolution

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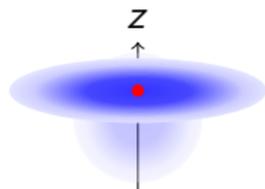
- low detection noise
($SNR \ll \sqrt{N}$)
- low information loss ($n_\gamma \ll 1$)

⇒ high resolution non-destructive detection

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CLASSICAL NON-DESTRUCTIVITY

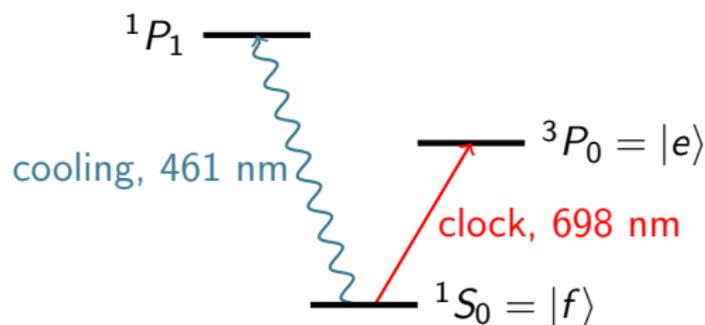
- Low photon scattering ⇒ atoms stay trapped
- Atoms recycle ⇒ less dead time in the clock cycle
⇒ reduced Dick effect

1 OPTICAL LATTICE CLOCKS

2 BEYOND THE QUANTUM PROJECTION NOISE

3 NON-DESTRUCTIVE DETECTION

MEASURING p WITH SR ATOMS

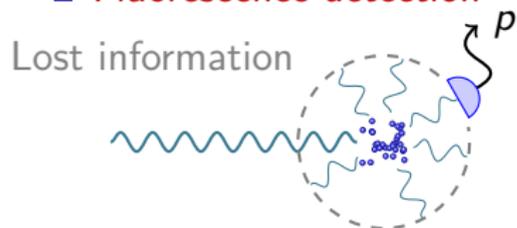


Probing the $^1S_0 \rightarrow ^1P_1$ transition: measure of $N_{|f\rangle}$

$$p = 1 - \frac{N_{\text{total}}}{N_{|f\rangle}}$$

USUAL SCHEME: FLUORESCENCE DETECTION

■ Fluorescence detection

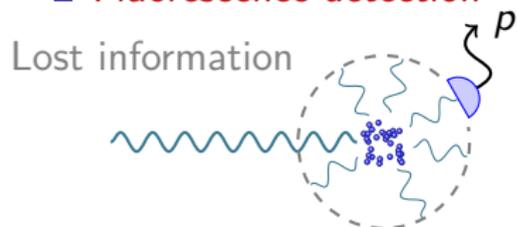


- Low efficiency \Rightarrow powerful probe beam
- Destructive detection: the atoms are scattered and lost ($n_\gamma \gg 1$)

DETECTION METHODS

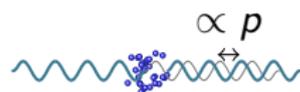
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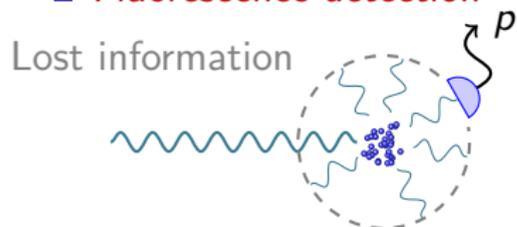


- Phase shift \Rightarrow low power probe beam
- SNR fundamentally limited by the light shot noise

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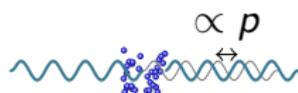
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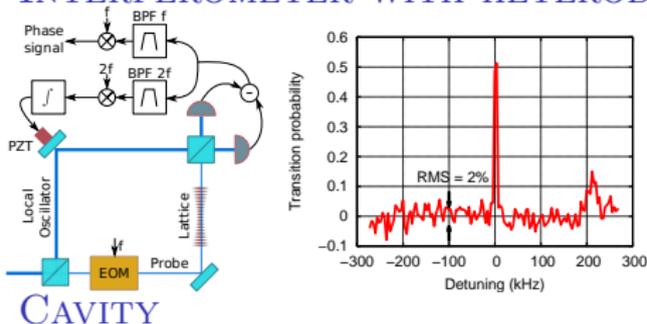
NON-DESTRUCTIVE DISPERSIVE DETECTION



- Phase shift \Rightarrow low power probe beam
- SNR fundamentally limited by the light shot noise
- Classical non-destructivity
- Quantum non-destructivity: no information loss

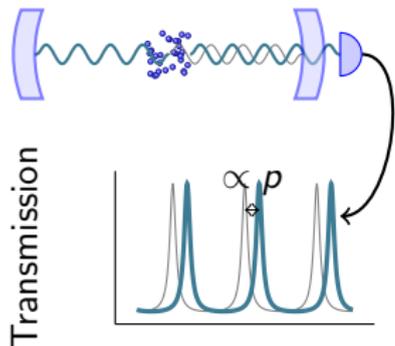
NON-DESTRUCTIVE DETECTION

INTERFEROMETER WITH HETERODYNE DETECTION



- Shot noise limited SNR
- Classical non-destructivity

J. Lodewyck et al. Phys. Rev. A **79** 061401(R) (2009)



- Longer interaction
 \Rightarrow improved SNR ($\times \sqrt{\mathcal{F}}$)
- Self-alignment
- Long term mechanical stability

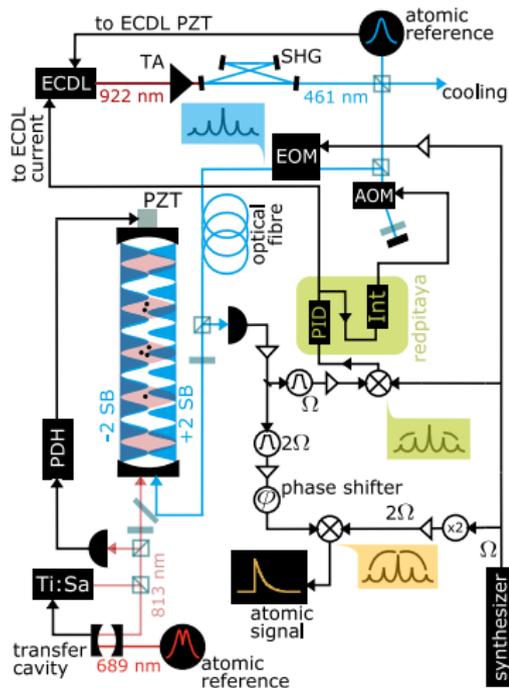
EXPERIMENTAL SETUP

CHALLENGE

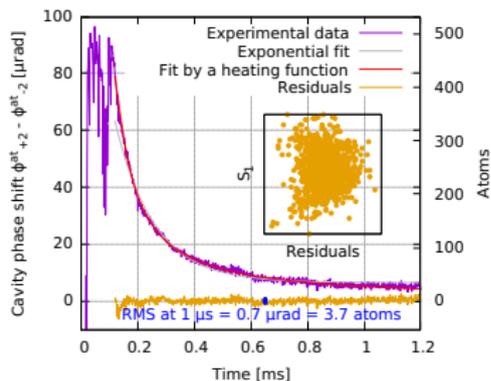
- discriminate **technical noises** from the atomic signal

DESIGN

- Bi-chromatic cavity 813 nm + 461 nm
⇒ lattice and detection aligned
- High finesse (16 000) at 461 nm
⇒ **100 fold increase of the SNR**
- Dual mode injection
⇒ **Immune to technical fluctuations** (cavity, laser)
⇒ **Homogeneous atom-cavity coupling**
- Heterodyne, PDH-like, detection ⇒ close to the **shot noise limit**

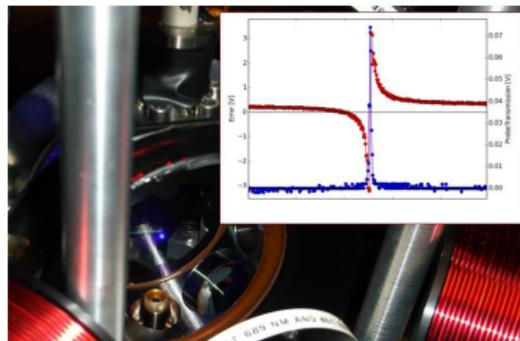


EXPERIMENTAL RESULTS

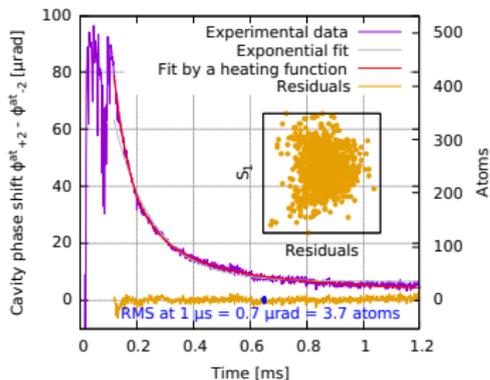


DETECTION SIGNAL

- Dynamic range of $\simeq 500$ atoms
- Scattering rate well modeled
- Immunity to technical noises demonstrated



EXPERIMENTAL RESULTS

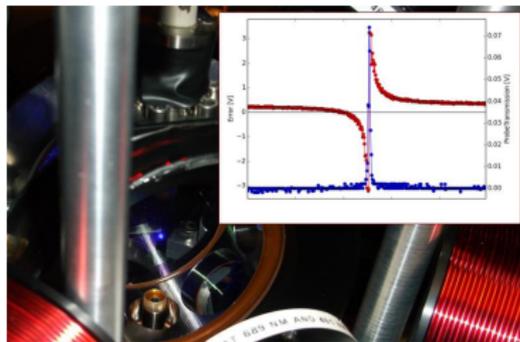


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DETECTION NOISE δN

- $\delta N = 23 \text{ atoms} / \sqrt{n_\gamma}$
- **Classical non-destructive regime**
 $\delta N = 3.7 \text{ atoms}$ for $n_\gamma = 38$ photons
 \Rightarrow high resolution
- **Quantum non-destructive regime**
 $\delta N > 23 \text{ atoms}$ for $n_\gamma < 1$ photon
 $\Rightarrow \delta N < \sqrt{N}$ for $N > 530$ atoms.



CONCLUSION AND PROSPECTS

REQUIREMENTS FOR A CLOCK DETECTION IN THE QUANTUM REGIME

- High SNR (cavity assisted)
- Low destructivity
- Homogeneous coupling
- Robust for operating in a state-of-the-art optical clock

PROSPECTS

- **Classical non-destructivity** for an improved frequency stability
- Demonstrate **quantum correlations**
(technical issues for a low scattering)
- **Overcome the QPN** limit for the frequency stability

Post-doc position available