

Novel source of multimode squeezed light for quantum enhanced space-time positioning

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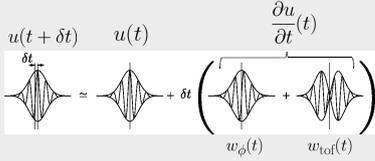
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We present a new multimode squeezed-states quantum source operating in a single pass regime suitable for ultimate precision space-time positioning experiments.

Space-time Positioning

Finding out position in space by ranging to a reference.
Measurement of time delay changes δt in a reflected pulse.



Enhancement of accuracy using the complete detection mode in homodyne detection [3]

Choice of detection mode

Coherent interferometric phase measurement

$$w_{LO} = w_{\phi}(t)$$

$$(\delta t)^{\text{phase}} = \frac{1}{2\sqrt{N}\omega_0}$$

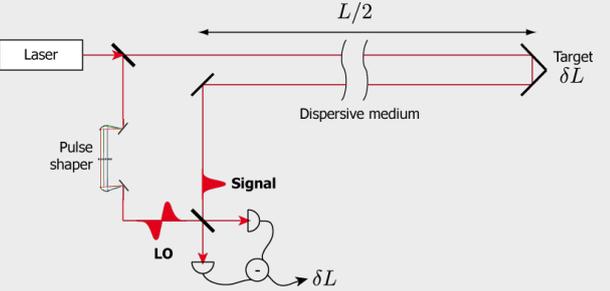
Incoherent time-of-flight measurement

$$w_{LO} = w_{\text{tof}}(t)$$

$$(\delta t)^{\text{tof}} = \frac{1}{2\sqrt{N}\Delta\omega}$$

Detection mode

$$w_L(\omega) = \frac{1}{\sqrt{\frac{\omega_0^2}{v_{\phi}^2} + \frac{\Delta\omega^2}{v_g^2}}} \left(\frac{\omega_0}{v_{\phi}} w_{\phi}(\omega) + \frac{\Delta\omega}{v_g} w_{\text{tof}} \right)$$

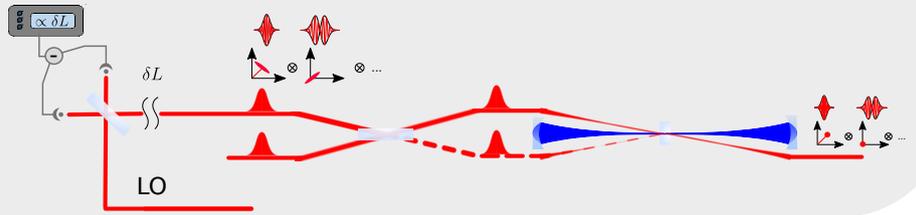


Minimum resolvable distance

$$\delta L_{\text{min}} = \frac{1}{2\sqrt{N}} \left(\left(\frac{\omega_0}{v_{\phi}} \right)^2 + \left(\frac{\Delta\omega}{v_g} \right)^2 \right)^{-1/2}$$

Beyond the Standard Quantum Limit

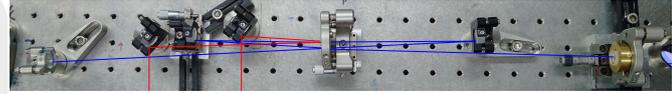
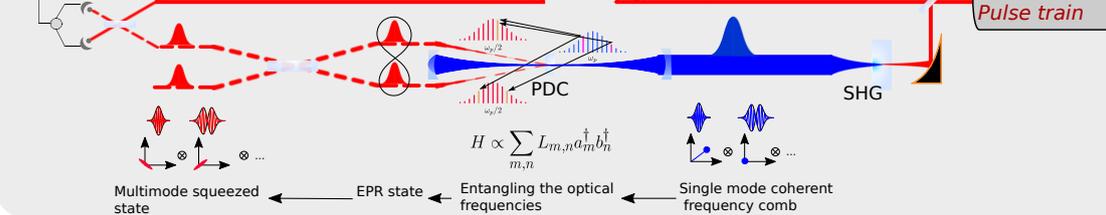
Further reduction of the minimum resolvable distance if the input beam has a mean field mode squeezed in the phase quadrature and a time-of-flight vacuum mode squeezed in the amplitude quadrature [4]



Experimental Setup

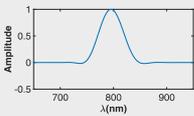
Fast Scope

Homodyne detection

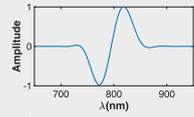


Our source consists in a synchronously pumped optical cavity resonant for the pump of a non-collinear type I parametric down-conversion process. This kind of cavity provides a power enhancement of the whole frequency comb which resonates in it. Due to its high peak power the resonating pulse is able to generate an EPR state which has a high multipartite entanglement between the frequency components of the signal and the idler pulses. When the EPR is sent to a beamsplitter a multimode squeezed state is obtained and can be analyzed by homodyne detection with a properly shaped local oscillator.

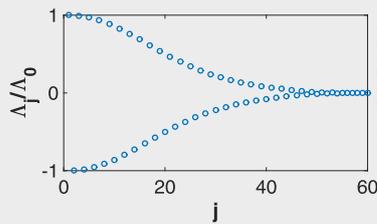
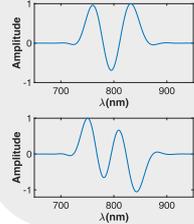
Multimode Quantum State Characterization



It is always possible to find a basis that diagonalize the interaction Hamiltonian describing the evolution of the multimode squeezed state. In this basis the quantum state is found to be a set of independently squeezed modes [1].



The spectral shape of these eigenmodes is found diagonalizing the phase matching matrix \mathcal{L}_{mn} whose eigenvalues are proportional to the squeezing parameter related to the corresponding eigenmode. The eigenmode basis is a set of Hermite-Gaussian polynomials.



$$\begin{pmatrix} \hat{A}_m \\ \hat{B}_m \end{pmatrix} = \begin{pmatrix} \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{pmatrix} \begin{pmatrix} \hat{a}_m \\ \hat{b}_m \end{pmatrix}$$

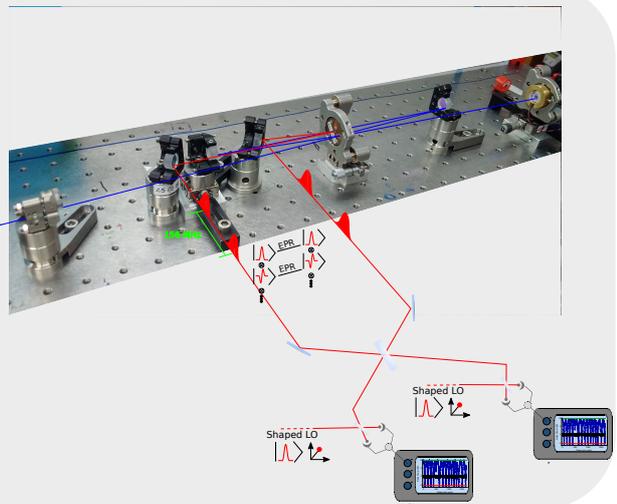
$$H = i\frac{\hbar}{2} k (\hat{A}^\dagger)^T \mathcal{L} \hat{A}^\dagger + \text{h.c.}$$

$$\mathcal{L}_{mn} = \text{sinc} \left[(k_{m+n}^{(p)} - k_m^{(s)} - k_n^{(i)}) \ell \right] \alpha_{m+n}$$

$$U \mathcal{L} U^T = \Lambda$$

$$\hat{S}^\dagger = U \hat{A}^\dagger \Rightarrow H = i\hbar \sum_j \Lambda_j (\hat{S}_j^\dagger)^2 + \text{h.c.}$$

Via pulse-shaping we can shape the local oscillator in a basis of frequency pixels. Different homodyne measurements then allow to build a covariance matrix. The diagonalization of this matrix gives the eigenmodes of the quantum state.



References:

- (1) J. Roslund, R. M. de Araujo, S. Jiang, C. Fabre, and N. Treps, *Nature Photonics*, **8**, 109112 (2014)
- (2) R. Raussendorf, H. J. Briegel, *Phys. Rev. Lett.*, **86**(22), 5188, (2001)
- (3) B. Lamine, C. Fabre, and N. Treps, *Phys. Rev. Lett.*, **101**, 123601 (2008)
- (4) O. Pinel, J. Fade, D. Braun, P. Jian, N. Treps and C. Fabre, *Phys. Rev. A*, **85**, 010101, (2011)