



## University of California, Berkeley

Holger Müller

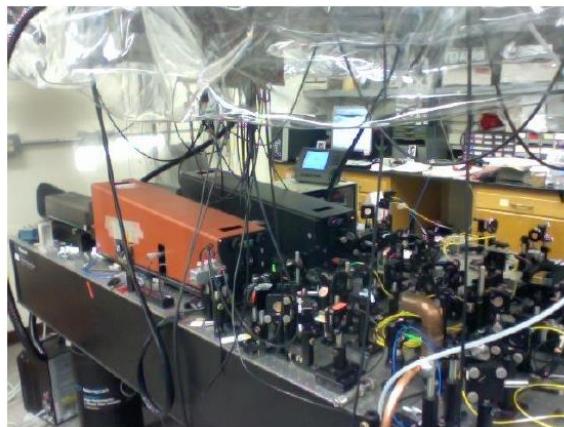
Brian Estey

Chenghui Yu

## Nanyang Technological University

Shau-Yu Lan

Pei-Chen Kuan



# Determination of Fine Structure Constant from Photon Recoil Frequency

Rydberg constant in terms of energy  $hcR_\infty = \frac{1}{2}m_e c^2 \alpha^2$

$$\alpha = \left[ \frac{2}{c} \frac{R_\infty}{m_e} \frac{M}{u} \frac{h}{M} \right]^{1/2}$$

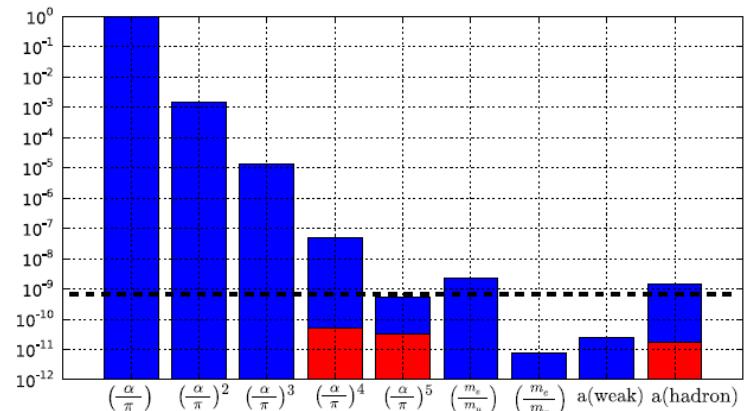
Rydberg Constant      Electron mass      Cs mass

Determined by the photon recoil frequency

$$\frac{h}{M} = \frac{4\pi c^2 \omega_r}{\omega^2}$$

LKB-11 :  $6.6 \times 10^{-10}$  PRL **106**, 080801 (2011)

Harvard-08 :  $3.7 \times 10^{-10}$  PRL **100**, 120801 (2008)



# Ramsey Interferometer

PHYSICAL REVIEW

VOLUME 78, NUMBER 6

JUNE 15, 1950

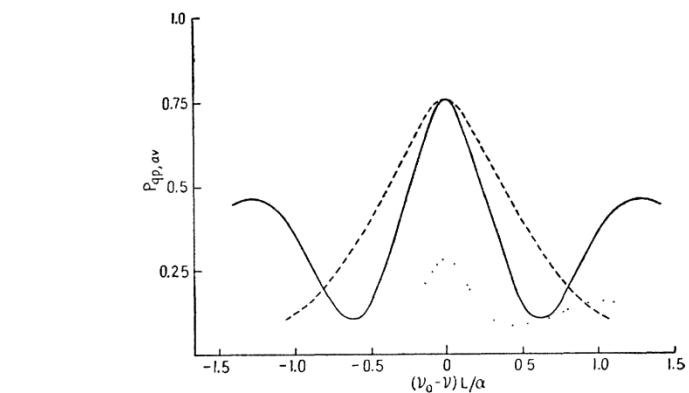
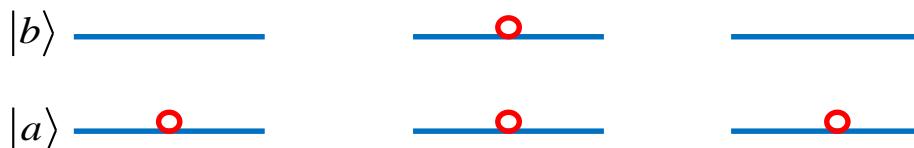
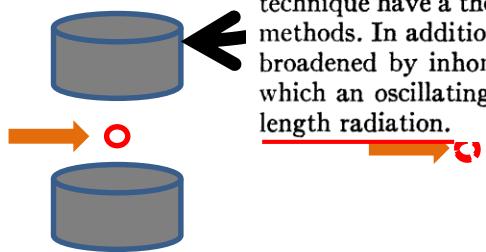
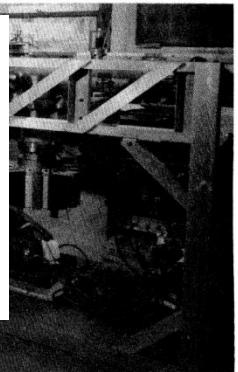
## A Molecular Beam Resonance Method with Separated Oscillating Fields\*

NORMAN F. RAMSEY

*Harvard University, Cambridge, Massachusetts*

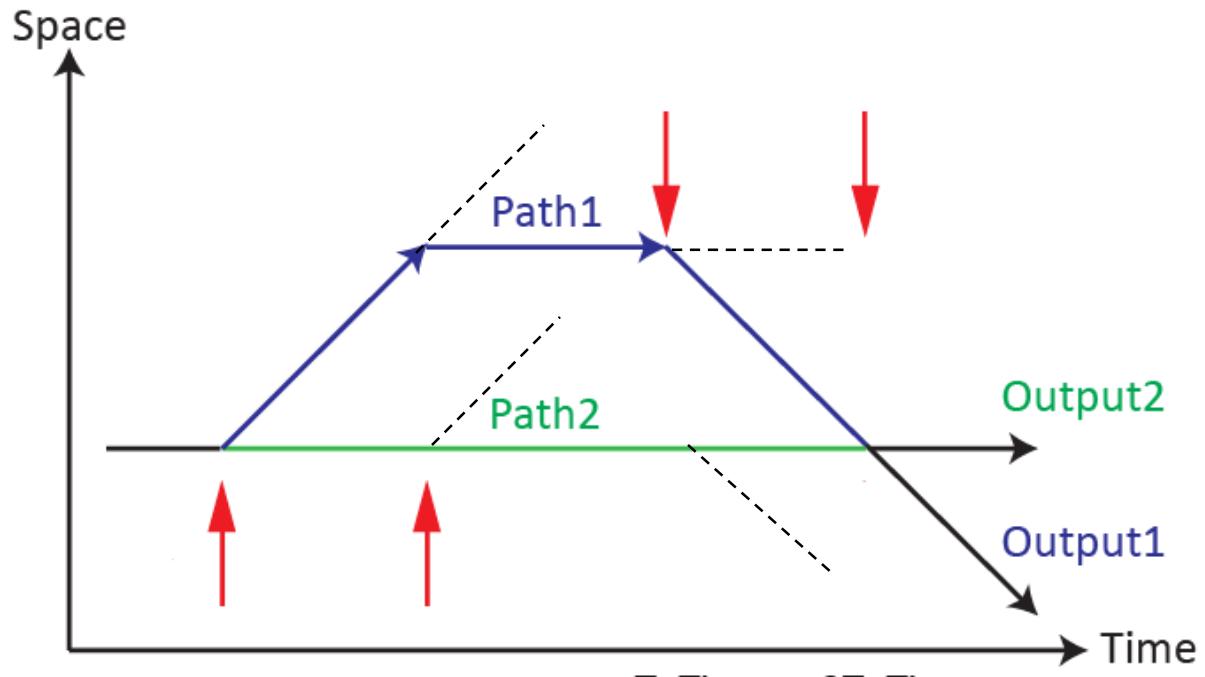
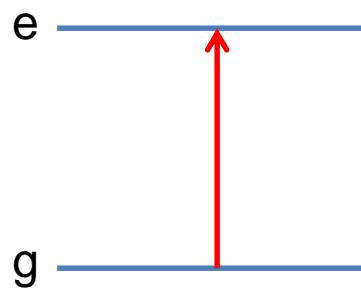
(Received February 21, 1950)

A new molecular beam resonance method using separated oscillating fields at the incident and emergent ends of the homogeneous field region is theoretically investigated in this paper. An expression is obtained for the quantum mechanical transition probability of a system between two states when the system is subjected to such separated oscillating fields. This is numerically averaged over the molecular velocity distribution and provides the theoretical shape of the resonance curves. It is found that resonances with such a technique have a theoretical half-width only 0.6 as great as those by conventional molecular beam resonance methods. In addition to producing sharper resonance minima, the new method has its resonances much less broadened by inhomogeneities of the fixed field, it makes possible resonance experiments in regions into which an oscillating field cannot be introduced, and it is more convenient and effective with short wavelength radiation.

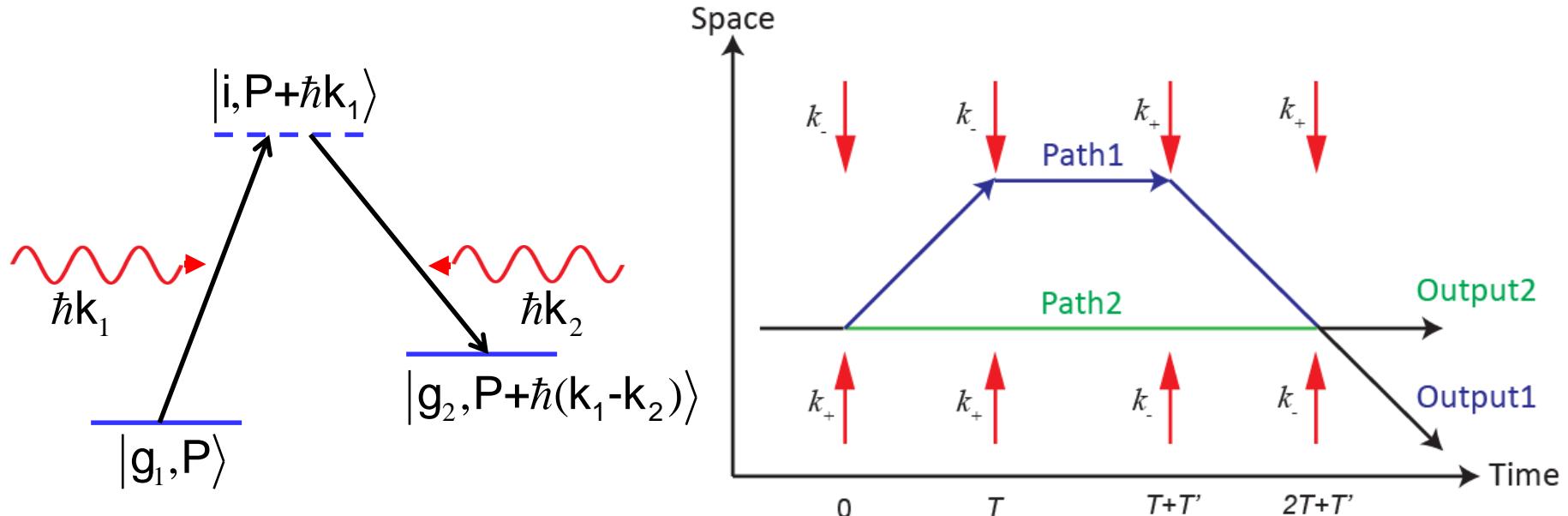


# Ramsey-Bordé Interferometer

- Designed to improve microwave clock
- Significant photon recoil shift  $\sim \text{mm/s}$
- Two additional pulses to close the interferometer



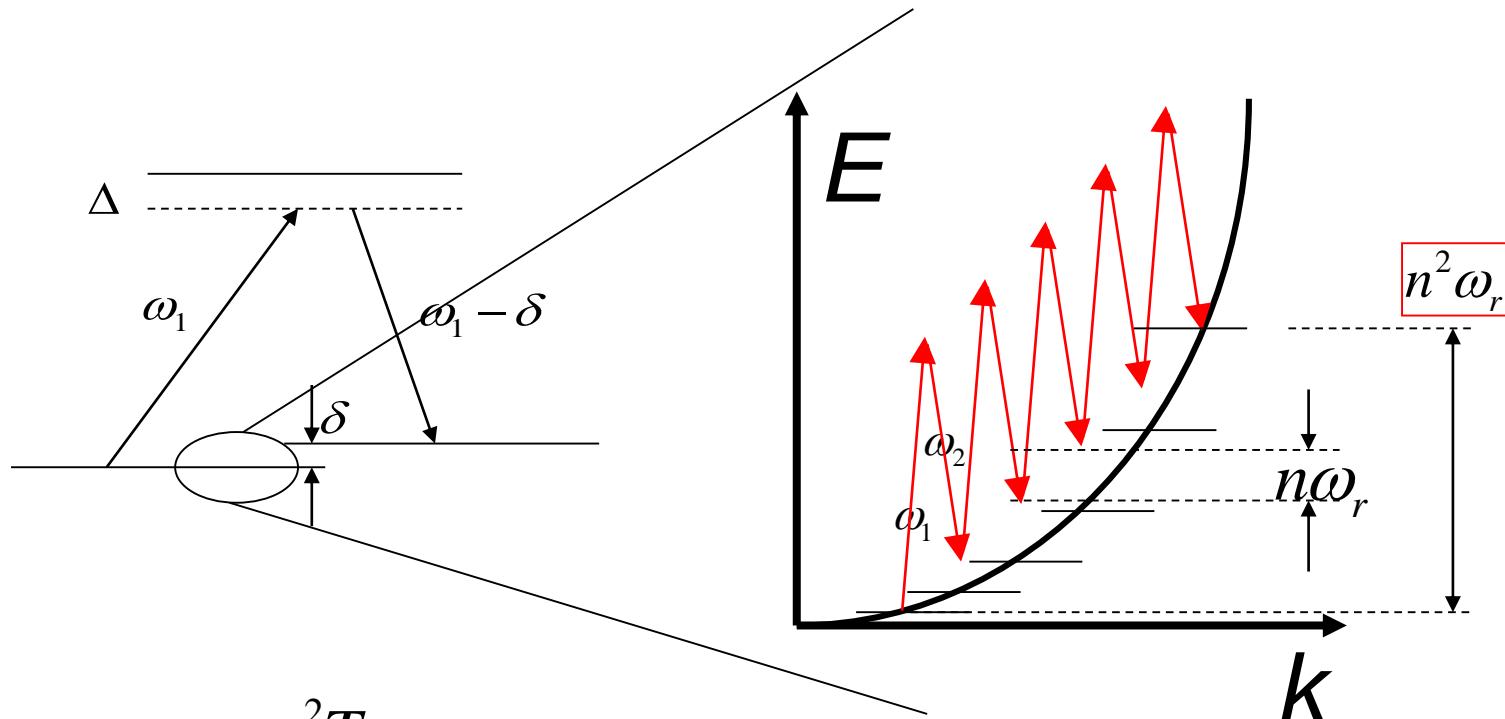
# Stimulated Raman Transition



$$\Delta\Phi = 8n^2 \omega_r T + n\Phi_L \quad \text{where } \omega_r = \frac{1}{2} \frac{\hbar}{m} k^2$$

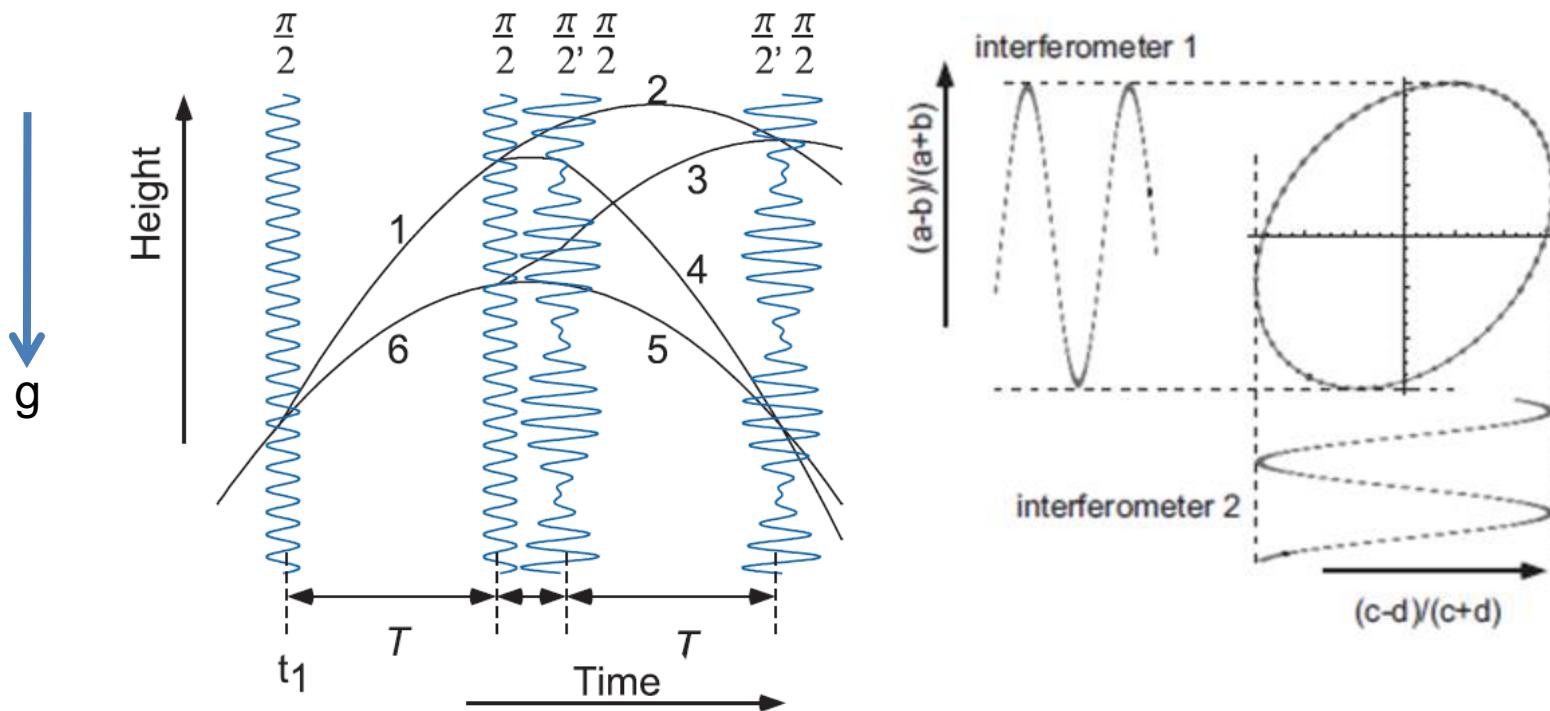
# Large Momentum Transfer

## Multi-Photon Bragg Diffraction



- Sensitivity  $\propto n^2 T$
  - $|g, p = p_0\rangle \rightarrow |g, p = n\hbar k_{eff}\rangle$
- $\Rightarrow$  external field insensitive

# Simultaneous Conjugate Interferometer



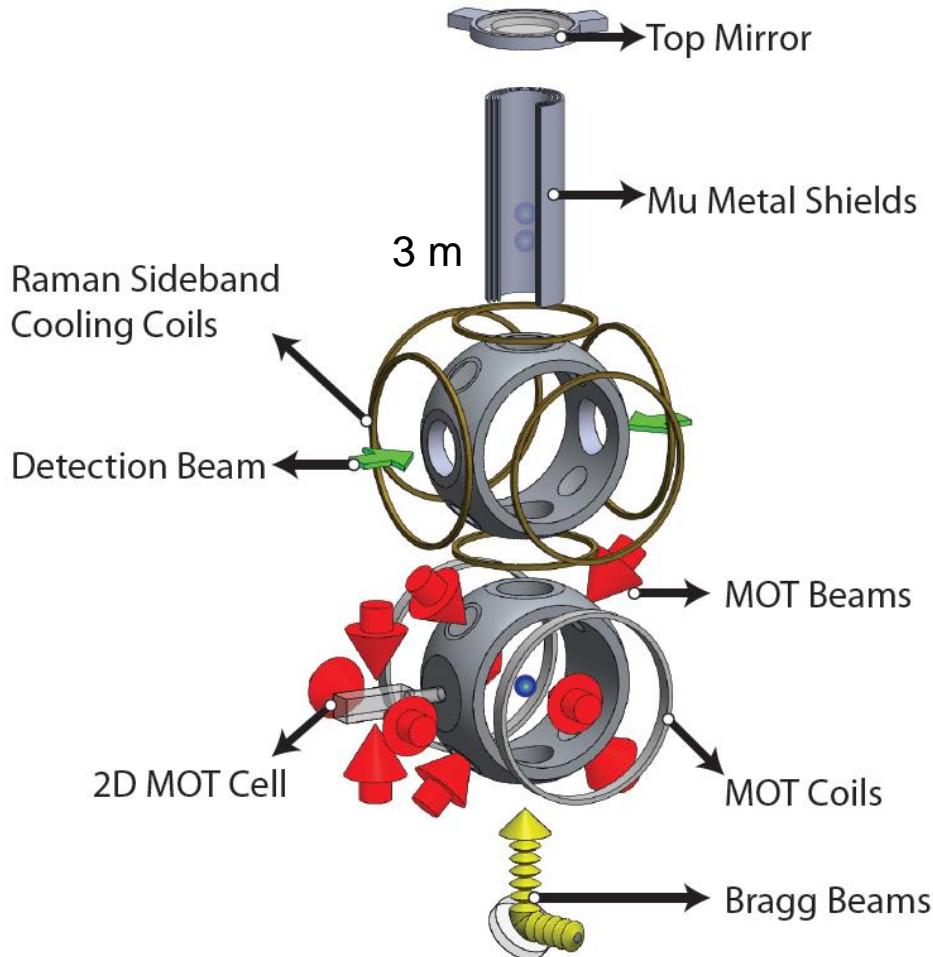
$$\Delta\Phi = \Delta\Phi^+ - \Delta\Phi^- = 16n^2\omega_r T + n\Phi_{\text{beamsplitter}}$$

$$\Phi_{\text{beamsplitter}} = (\Phi_3^+ - \Phi_3^-) + (\Phi_4^- - \Phi_4^+)$$

Common Mode Rejection

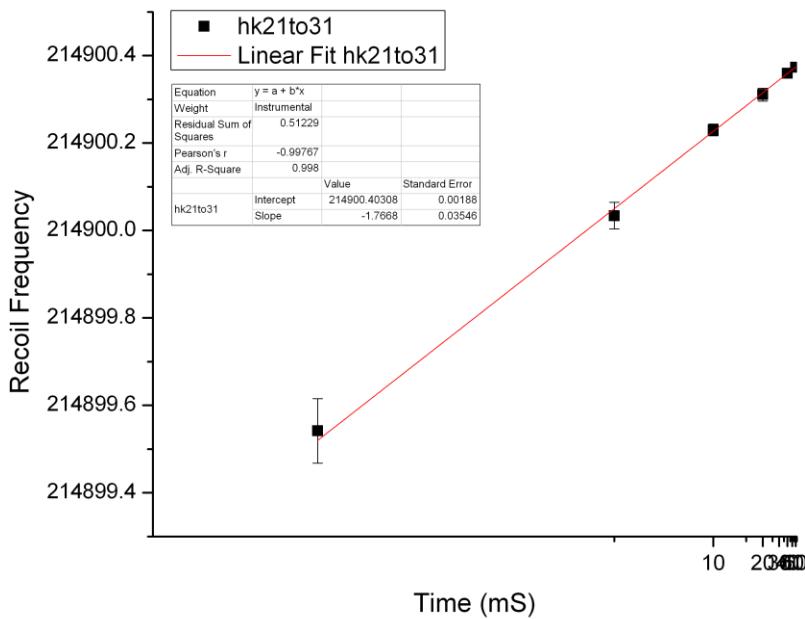
# Experimental Setup

## Atomic Fountain

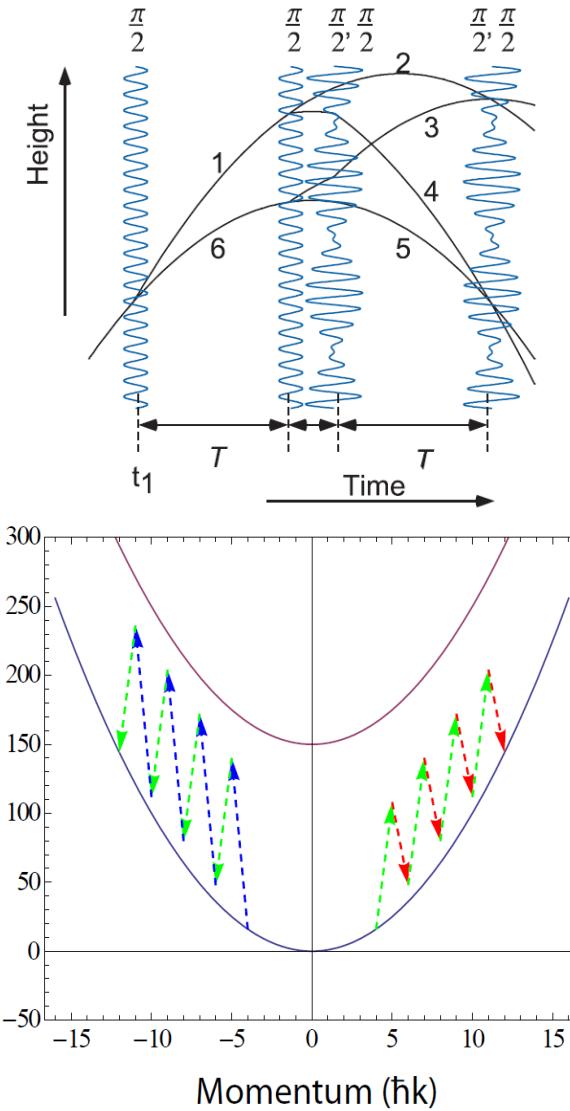


- 3D magneto-optical trap
- Molasses cooling
- Adiabatic cooling
- Launch
- Raman sideband cooling
- Adiabatic microwave transfer
- $F=3, m=0$  state selection
- Velocity selection
- Interferometer sequence
- Fluorescence detection

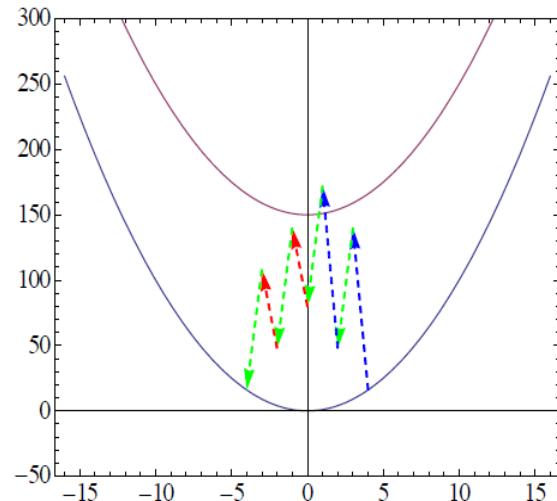
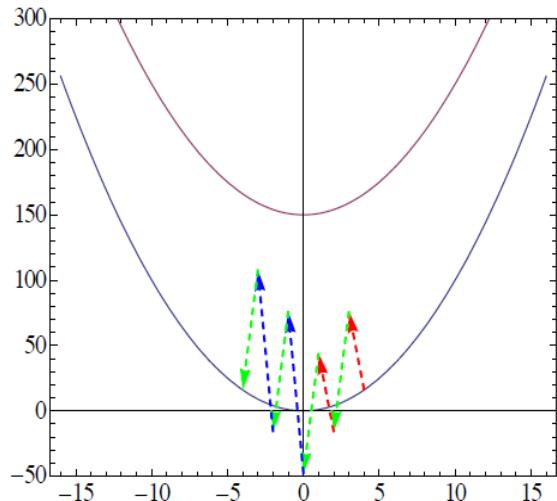
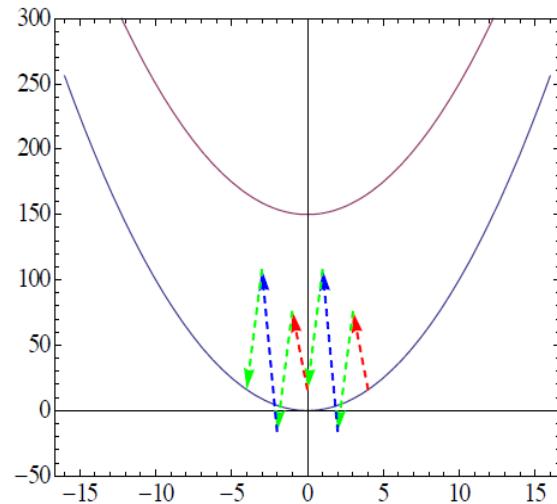
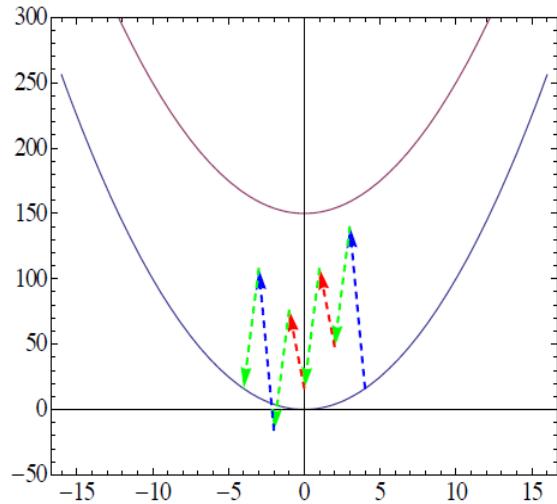
# Diffraction phase shift



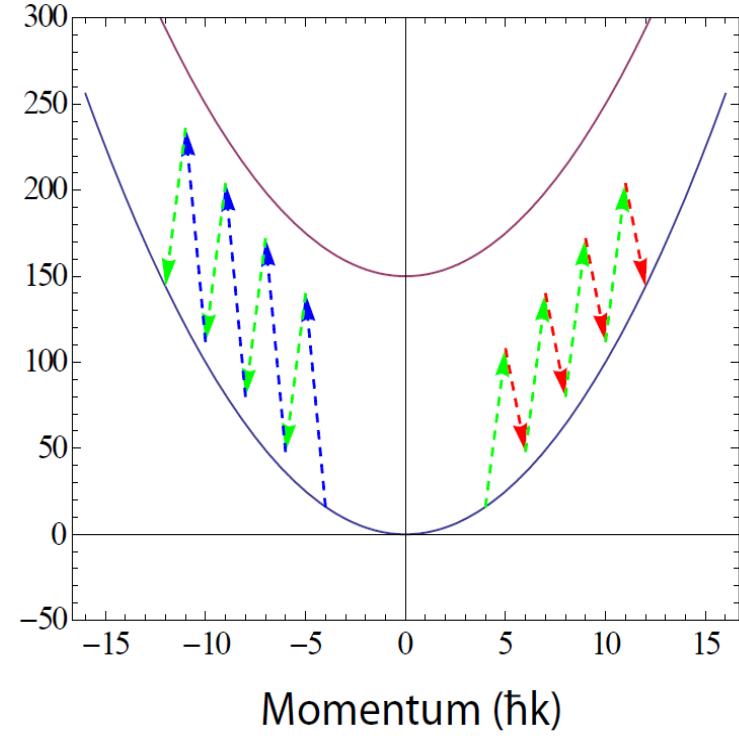
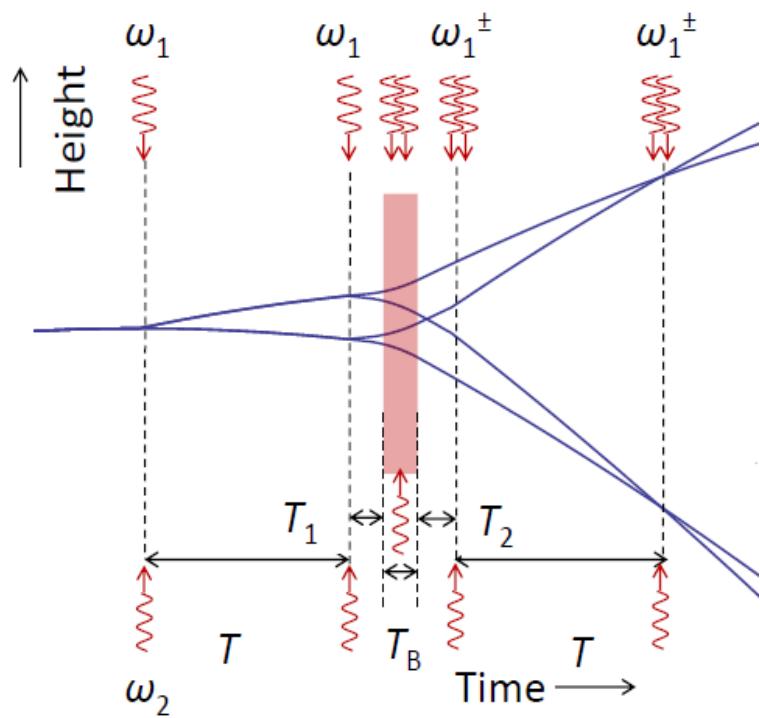
$$\omega_m = \omega_r + \frac{\varphi_o}{T}$$



# Cross coupling

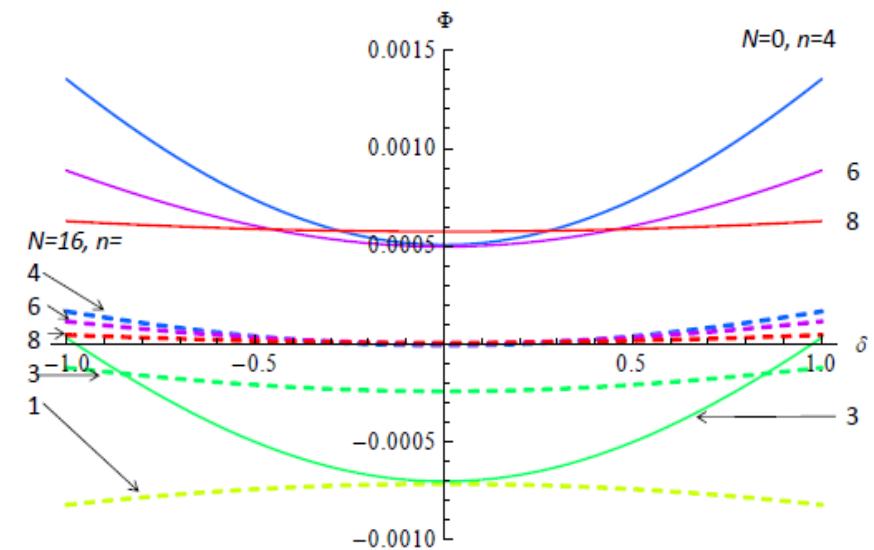
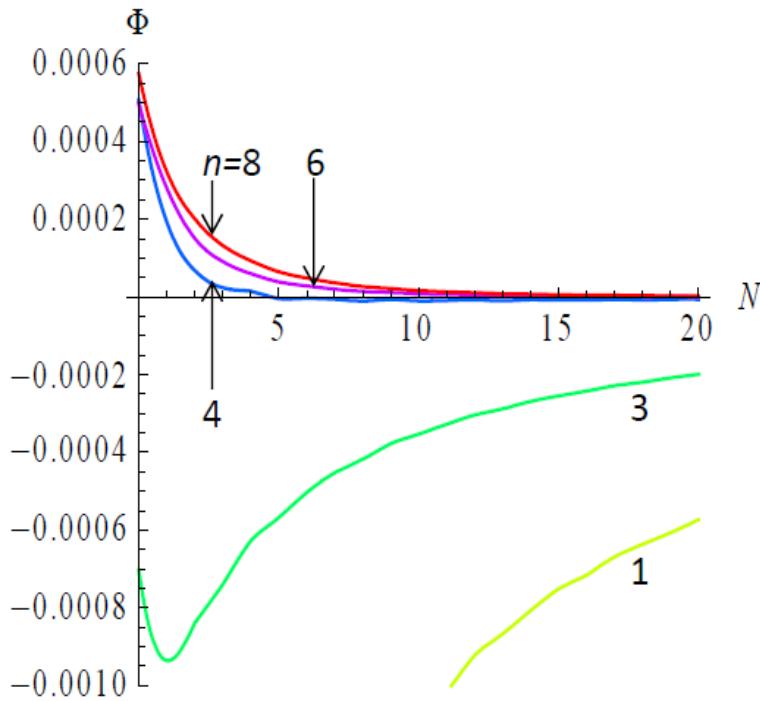


# Bloch Oscillation

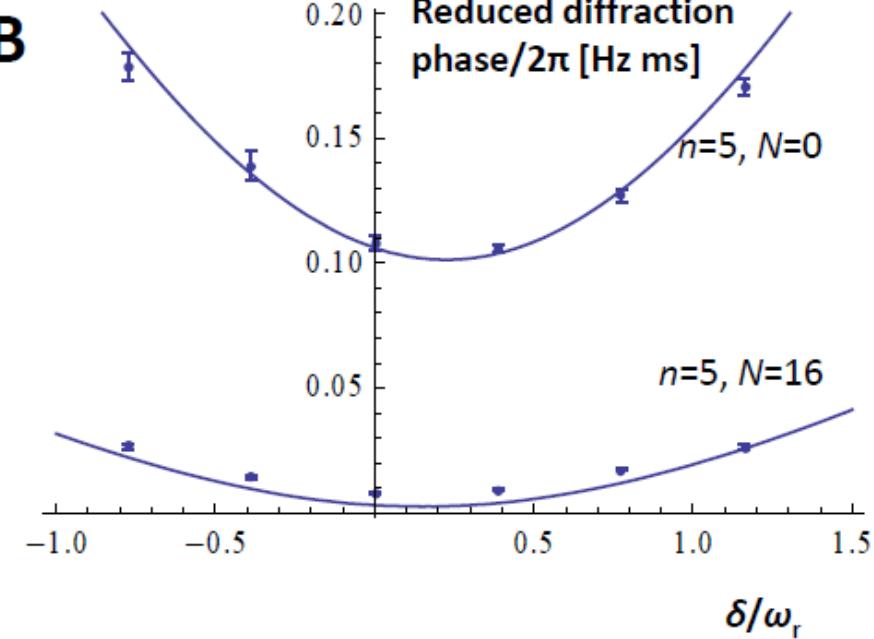
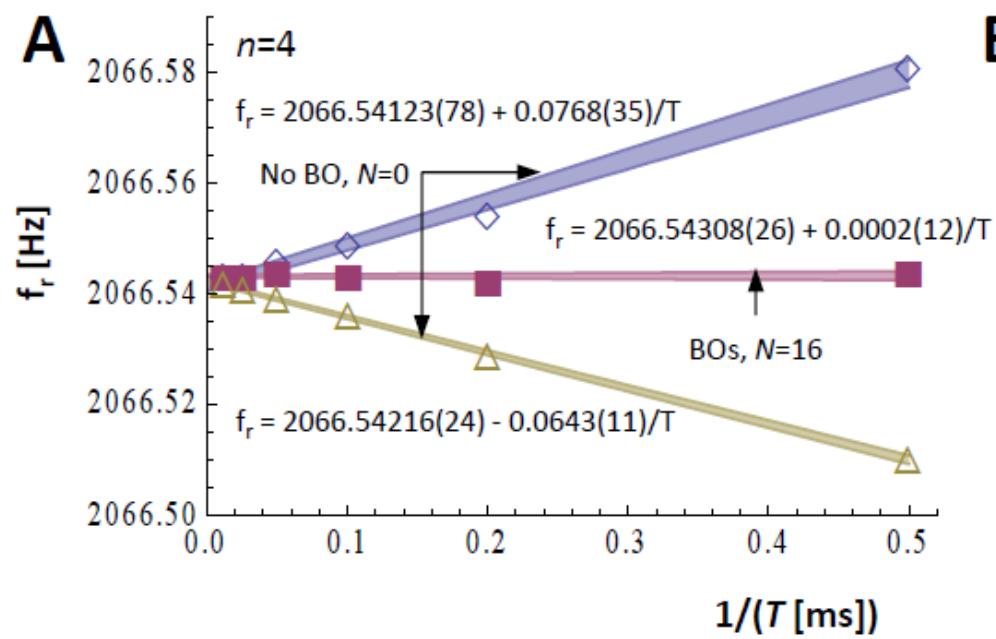


$$\Phi = 16n(n + N)\omega_r T$$

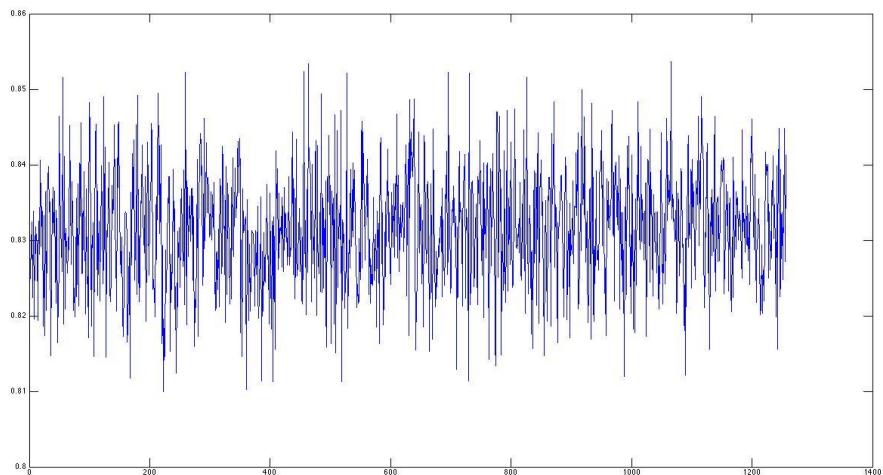
# Simulation of diffraction phase



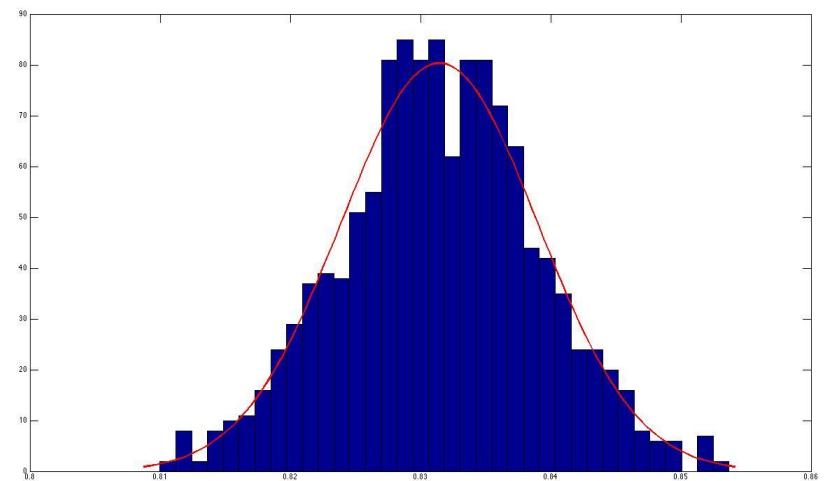
# Suppression of diffraction phase



# Long data



0.33 ppb/sqrt(6hr)



# More Systematic Effects

- Laser frequency 351,730,902,093(10) kHz
- Gravity gradient -2.954(85) ppb
- Guoy phase -1.902(46) ppb
- Wavefront curvature 0.111(55) ppb
- Beam alignment -0.085(49) ppb
- Zeeman effect -0.051(14) ppb
- Mean Field Shift
- Value Different Bragg order
- Atom temperature effect
- Detection method
- 
- 
-



# Thank You

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